## EXAM QUANTUM THEORY, 13 JANUARY 2020, 10.15-13.15 HOURS.

- 1. The wave function of a particle in position representation is  $\psi(x)$ , normalized to unity:  $\int_{-\infty}^{\infty} |\psi(x)|^2 dx = 1$ . (The particle is spinless and confined to the *x*-axis.)
- *a*) What is the corresponding wave function  $\psi(p)$  in momentum representation? Verify that it is also normalized to unity.
- *b*) The time-reversal operation *T* in position representation is
   *T*ψ(x) = ψ\*(x). What is the corresponding time-reversal operation in momentum representation?
- *c)* The Kramers theorem says that, under certain conditions, the eigenstates in the presence of time reversal symmetry are twofold degenerate. Does Kramers theorem apply in this case? Motivate your answer.
- 2. Consider the operator  $T_a = e^{iap/\hbar}$ , where *a* is a real constant and  $p = -i\hbar d/dx$  is the momentum operator along the *x*-axis.
- *a*) Explain why  $T_a$  is the operator for translation over a distance *a*.

An electron in a chain of atoms at positions  $x_n = na$ ,  $n = 0, \pm 1, \pm 2, ...$ , has wave function amplitude  $\psi_n$  on atom n. The electron can move along the chain by hopping from one atom to the next.

• *b)* Explain why the Hamiltonian

 $H = \alpha (T_a + T_a^{\dagger})$ 

describes the electron motion for a real constant  $\alpha$ . How should *H* look like if  $\alpha$  was complex?

- *c)* Calculate the energy spectrum of *H*. (You may still assume a real *α*.) How does the expectation value *ν* of the velocity of the electron depend on its momentum *p*?
- 3. In this problem we will investigate a two-level system (ground state  $|g\rangle$  at energy  $-\epsilon/2$ , excited state  $|e\rangle$  at energy  $\epsilon/2$ ) interacting with a harmonic oscillator (frequency  $\omega$ , bosonic creation operator  $a^{\dagger}$ ). The Hamiltonian is

$$H = -\frac{\varepsilon}{2}|g\rangle\langle g| + \frac{\varepsilon}{2}|e\rangle\langle e| + \hbar\omega(a^{\dagger}a + \frac{1}{2}) + \gamma(|e\rangle\langle g|a + |g\rangle\langle e|a^{\dagger}),$$

where  $\gamma$  is the interaction energy. This socalled Jaynes-Cummings model describes how an atom (modelled by the two-level system) can be excited by absorbing a photon in an optical cavity (modelled by the harmonic oscillator).

• *a*) Show that the operator

$$N = a^{\dagger}a + \frac{1}{2}|e\rangle\langle e| - \frac{1}{2}|g\rangle\langle g|$$

commutes with *H*. What is the corresponding conserved quantity?

Because *N* is conserved, we can assume it has a definite value  $N_0$  and restrict the Hilbert space to just two states: a state  $|\psi_1\rangle = |g, N_0\rangle$  where the atom is in the ground state and the cavity contains  $N_0$  photons, and a state  $|\psi_2\rangle = |e, N_0 - 1\rangle$  where the atom is in the excited state and the cavity contains  $N_0 - 1$  photons.

• *b*) Calculate the four matrix elements  $\langle \psi_1 | H | \psi_1 \rangle$ ,  $\langle \psi_1 | H | \psi_2 \rangle$ ,  $\langle \psi_2 | H | \psi_1 \rangle$ ,  $\langle \psi_2 | H | \psi_2 \rangle$ .

*Help:* Recall that the harmonic oscillator at energy  $(n + \frac{1}{2})\hbar\omega$  is in the state  $|n\rangle = (n!)^{-1/2}(a^{\dagger})^n |0\rangle$ .

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The four matrix elements can be combined into the  $2 \times 2$  matrix

$$M = \text{constant} \times I + \begin{pmatrix} -\delta/2 & \gamma\sqrt{N_0} \\ \gamma\sqrt{N_0} & \delta/2 \end{pmatrix},$$

where *I* is the unit matrix and  $\delta = \varepsilon - \hbar \omega$ .

- *c)* The two eigenstates of *H* at a given  $N_0$  have an energy difference  $\delta E$ . Calculate  $\delta E$  and check that  $\delta E \rightarrow \varepsilon \hbar \omega$  in the limit  $\gamma \rightarrow 0$ .
- 4. In class we studied the Casimir effect with two metal plates, here we consider instead three parallel metal plates, at  $x = L_1$ , at  $x = L_2$ , and at  $x = L_3$ , with  $0 < L_1 < L_2 < L_3$ . The plates have infinite extent in the y-z plane. We treat the electromagnetic field as a scalar field  $\phi(x, y, z)$  (ignoring polarization). In vacuum the quantum fluctuations result in an energy per unit area given by  $E_{\text{total}} = E(L_2 L_1) + E(L_3 L_2)$ , where the function E(L) is given by

$$E(L) = \frac{1}{2}\hbar c \sum_{n=1}^{\infty} \int_{-\infty}^{\infty} \frac{dk_y}{2\pi} \int_{-\infty}^{\infty} \frac{dk_z}{2\pi} \sqrt{\pi^2 n^2 / L^2 + k_y^2 + k_z^2}$$
(1a)

$$=\frac{\hbar c}{4\pi}\sum_{n=1}^{\infty}\int_{\pi n/L}^{\infty}u^2\,du.$$
(1b)

• *a*) Explain how to arrive at each of these two equations (1a) and (1b) for E(L).

Because the integral over u diverges, we insert a factor  $e^{-\epsilon u}$  with  $\epsilon > 0$  and calculate as follows:

$$E(L,\epsilon) = \frac{\hbar c}{4\pi} \sum_{n=1}^{\infty} \int_{\pi n/L}^{\infty} e^{-\epsilon u} u^2 du$$
(2a)

$$=\frac{\hbar c}{4\pi}\frac{d^2}{d\epsilon^2}\frac{1}{\epsilon}\sum_{n=1}^{\infty}e^{-\epsilon\pi n/L}$$
(2b)

$$=\frac{\hbar c}{4\pi}\frac{d^2}{d\epsilon^2}\frac{1}{\epsilon}\frac{1}{e^{\epsilon\pi/L}-1}.$$
(2c)

• *b*) Explain the two steps in this calculation, from Eq. (2a) to (2b) and from Eq. (2b) to (2c).

We now take an infinitesimally small  $\epsilon$  and expand Eq. (2c) to arrive at

$$E(L,\epsilon) = \frac{\hbar c}{4\pi} \left( \frac{6L}{\pi\epsilon^4} - \frac{1}{\epsilon^3} - \frac{\pi^3}{360L^3} + \operatorname{order}(\epsilon^2) \right).$$
(3)

• *c*) Calculate the derivative  $F = -dE_{\text{total}}/dL_2$  the limit  $\epsilon \to 0, L_3 \to \infty$ . This is the force per unit area which attracts the second plate to the first plate, separated by a distance  $L_2 - L_1$ .