## Answers to the Exam Quantum Theory, RETAKE, 12 JANUARY 2014 each item gives 2 points for a fully correct answer, grade = total $\times 9/24 + 1$

- 1. a)  $\langle \psi | \rho \psi \rangle = \sum_{n} p_{n} |\langle \psi | \psi_{n} \rangle|^{2} \ge 0.$ b)  $d\rho/dt = \sum_{n} p_{n} (|d\Psi_{n}/dt\rangle \langle \Psi_{n}| + |\Psi_{n}\rangle \langle d\Psi_{n}/dt|) = (-i/\hbar) \sum_{n} p_{n} (H|\Psi_{n}\rangle \langle \Psi_{n}| - |\Psi_{n}\rangle \langle \Psi_{n}|H) = (-i/\hbar) [H, \rho].$ c)  $\rho(t) = e^{-iHt/\hbar} \rho(0) e^{iHt/\hbar}$  so  $\rho^{2}(t) - \rho(t) = e^{-iHt/\hbar} [\rho^{2}(0) - \rho(0)] e^{iHt/\hbar}$  and  $\rho^{2}(0) - \rho(0) = e^{iHt/\hbar} [\rho^{2}(t) - \rho(t)] e^{-iHt/\hbar}$ ; hence  $\rho^{2}(t) = \rho(t) \Leftrightarrow \rho^{2}(0) = \rho(0).$
- 2. *a*) *dE<sub>n</sub>/dλ* = (*d*/*dλ*)⟨*n*, *λ*|*H*(*λ*)|*n*, *λ*⟩ = ⟨*n*, *λ*|∂*H*/∂*λ*|*n*, *λ*⟩ plus a term consisting of the integral ∫[(*Hψ*)\*∂*ψ*/∂*λ* + (∂*ψ*\*/∂*λ*)*Hψ*]*dx*, with *ψ* the eigenstate of *H*. Because *Hψ* = *E<sub>n</sub>*(*λ*)*ψ*, this integral can also be written as *E<sub>n</sub>*(*λ*) ∫[*ψ*\*∂*ψ*/∂*λ* + (∂*ψ*\*/∂*λ*)*ψ*]*dx* = *E<sub>n</sub>*(*λ*)(*d*/*dλ*) ∫|*ψ*|<sup>2</sup>*dx* = 0 because of the normalization of *ψ*. *b*) the operator *p<sub>z</sub>* = -*iħ*∂/∂*z* commutes with *H*, so the Heisenberg equation of motion gives *dp<sub>z</sub>*/*dt* = 0. The velocity operator *v<sub>z</sub>* = (*p<sub>z</sub> eBy*)/*m* does not commute with *H*, so the velocity along the wire is not conserved. *c*) the velocity operator in the *z*-direction is *y* = ∂*H*/∂*n* : now use the

*c)* the velocity operator in the *z*-direction is  $v_z = \partial H / \partial p_z$ ; now use the Hellman-Feynman theorem,  $\langle v_z \rangle = dE(p_z)/dp_z$ .

- 3. *a*) the flux is  $\Phi = \pi |B|R_c^2 = \pi |B|(2mE)/(qB)^2$ , so quantization gives  $E_n = \hbar(|qB|/m)(n + \frac{1}{2})$ . The ground state is  $E_0 = |qB| \times (\hbar/2m)$ . *b*) In graphene  $E_n = \text{constant} \times \sqrt{n|B|}$ , so it increases more slowly with |B| and is zero for n = 0. *c*) motion along *z* is independent of motion in *x*-*y* plane, adds  $p_z^2/2m = (\hbar k)^2/2m$  to the energy;  $E_n(k) = (n + 1/2)\hbar |qB|/m + \hbar^2 k^2/2m$ , so spacing is  $\hbar |qB|/m$ .
- 4. *a*) The trace of H(t) is zero, so the eigenvalues are  $\pm E$ , and the determinant is  $-(\hbar e/2m)^2 B_0^2 = -E^2$ , so the eigenvalues are  $\pm \hbar e B_0/2m = \pm \frac{1}{2}\hbar\omega_0$ . *b*) The state  $\psi(0)$  is an eigenstate of H(0) with eigenvalue  $E = -\frac{1}{2}\hbar\omega_0$ , so the dynamical phase is  $c_1T = -iET/\hbar$  with  $c_1 = \frac{1}{2}\omega_0$ . *c*) The normalized and single-valued eigenstate of H(t) with eigenvalue  $E = -\frac{1}{2}\hbar\omega_0$  is  $|t\rangle = 2^{-1/2}(1, e^{2\pi i t/T})$ . The Berry phase is given by  $c_0 = i \int_0^T dt \langle t | \frac{d}{dt} | t \rangle = i \int_0^T dt \frac{1}{2}(2\pi i/T) {\binom{1}{e^{-2\pi i t/T}}} \cdot {\binom{0}{e^{2\pi i t/T}}} = -\pi$ .